# Non-equilibrium Thermodynamics

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#### The Equilibrium Picture

Fluctuations are small for macroscopic systems. (Large Numbers)

$$\frac{\Delta E}{\langle E \rangle} \simeq \frac{\sqrt{N}}{N} \simeq \frac{1}{\sqrt{N}}$$

System undergoes quasi-static changes and the net entropy production is 0. (Slow Dynamics, reversible)

Homogeneous systems can be described globally (Homogeneity)

Therefore, state variables are well defined!

#### The Equilibrium Picture

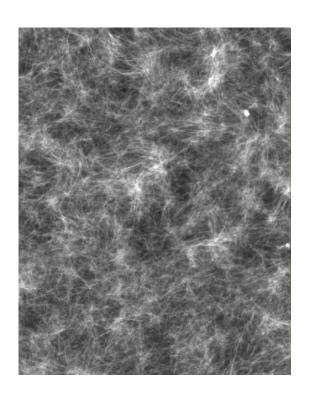
First Law:  $dU = \delta Q - \delta W$ 

Second Law: 
$$dS \ge \frac{\delta Q}{T}$$

## Nature is essentially non-equilibrium

#### **Active Matter**

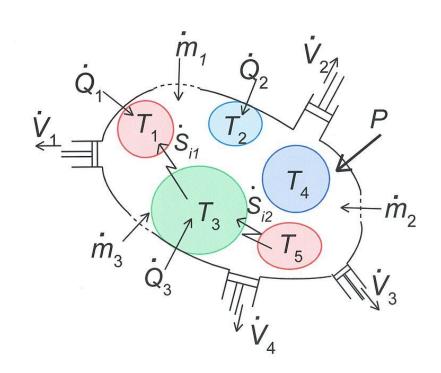




# Baby Step: Local Equilibrium in Non-equilibrium Systems (Near Equilibrium)

For inhomogeneous systems, state variables such as S,T,E are not well defined.

Instead, we invoke the concept of *local* equilibrium.



### **Entropy Production**

Locally, the generalized first law applies

$$\mathrm{d}u = T\mathrm{d}s + \sum_{i} \mu_{j} \mathrm{d}N_{j} + \phi \mathrm{d}\rho,$$
 Then, 
$$\mathrm{d}s = \frac{1}{T}\mathrm{d}u - \sum_{i} \left(\frac{\mu_{j}}{T}\right) \mathrm{d}N_{j} - \frac{\phi}{T}\mathrm{d}\rho_{\mathrm{e}},$$

In general,  $ds = \sum_k \phi_k d\rho_k$ , where phi is general potential.

	$\rho_k$	$\phi_{k}$	$ abla \phi_k$
energy	u	1/T	$\nabla(1/T)$
number of particles of type $j$	$N_j$	$\mu_j/T$	$\nabla(\mu_j/T)$
charge density	$ ho_{ m e}$	$-\phi_{ m e}/T$	$-\nabla(\phi_{\rm e}/T)$

#### **Entropy Production**

We are interested in the local entropy production rate

$$\Sigma = \frac{\partial s}{\partial t} + \nabla \cdot \boldsymbol{J}_s.$$

Where J\_s is the entropy current density, which can be calculated with

$$\boldsymbol{J}_s = \sum_k \phi_k \boldsymbol{J}_k.$$

Therefore, eventually

$$\Sigma = \sum_{k} \nabla \phi_k \cdot \boldsymbol{J}_k.$$

#### **Kinetic Coefficients**

So how do we know the generalized current density J\_k?

Assuming a linear response, we can relate them to the generalized potential by

$$\boldsymbol{J}_i = \sum_j L_{ij} \nabla \phi_j,$$

Then 
$$\Sigma = \sum_{ij} \nabla \phi_i L_{ij} \nabla \phi_j.$$

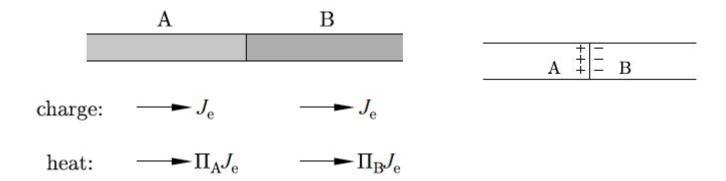
#### **Kinetic Coefficients**

Remarkably, we can show that the kinetic coefficients in  $\Sigma = \sum_{ij} \nabla \phi_i \, L_{ij} \, \nabla \phi_j.$  satisfies

$$L_{ij} = L_{ji}$$
.

This is called Onsager's reciprocal relations and can be proved for system in *near* equilibrium states, which means small perturbation and linear response.

## Application: Thermoelectricity



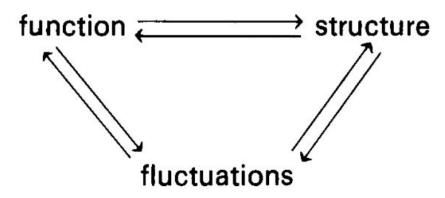
# Far Away from Equilibrium?

Drosophila embryonic development

#### Dissipative Structures

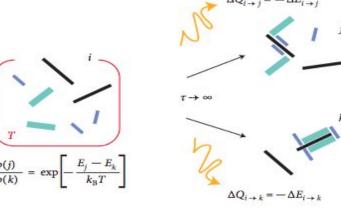
Local equilibrium is not satisfied for systems so far away from equilibrium!

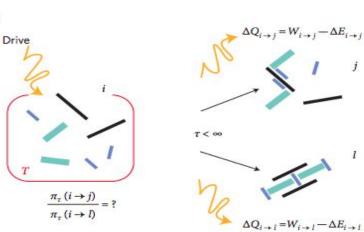
Ilya Prigogine proposed instead that we can describe these exotic systems as dissipative structures.



Application: Dissipative adaptation in driven self-assembly

Jeremy England, *Nature Nanotechnology* 2015





#### Conclusion

Global state variables break down for inhomogeneous systems at non-equilibrium.

At near equilibrium we can define local equilibrium and study the local entropy production through the generalized potentials.

However, far away from equilibrium, these concepts all break down. Surprisingly, new order emerges from the large fluctuations described by dissipative structures.

Non-equilibrium thermodynamics is at its infancy and an exciting field for young people (you!) to explore.