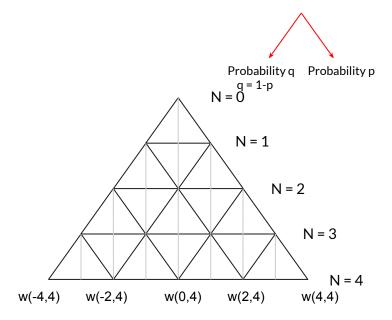
Diffusion and Random Walks

Emma Suen-Lewis

1-D Random Walk

Generally characterized by: time step Δx , spacial step Δt , probability of going left or right

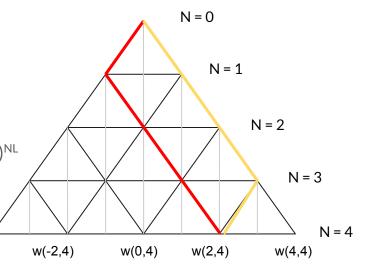
- Let p be chance of going right, q = 1- p be the chance of going left
- Let m = $x/\Delta x$ be cumulative number of spacial steps
- Let N = $t/\Delta t$ be variable for time steps passed
- Let w(m,N) be the probability of being at position m at time
 N



Finding the probability w(m,N) for a 1-D walk

W(-4,4)

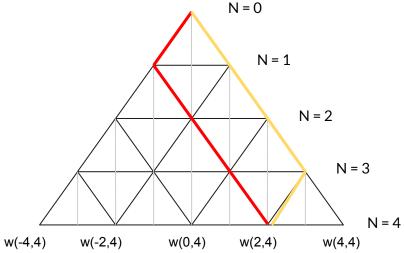
- For some N, only certain m are possible
 - o I.e. at N = 2, m is either -2, 0, or 2
- For some m and N, there is only 1 combination of right/left paths that lead to that point:
 - $\circ \qquad m = N_R N_L \text{ and } N = N_R + N_L$
 - \circ N_R = (N+m)/2 and N_L = (N-m)/2
- Since we know the probability of each step, for some path there is a probability that exactly that path is taken: p^{NR}(1-p)^{NL}



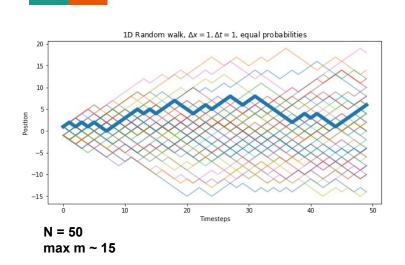
Finding the probability w(m,N) for a 1-D walk

• To get total probability for all paths leading to some m, N, multiply by the number of possible paths:

$$w(m,N) = p^{NR}(1-p)^{NL*} (N!/N_R!N_L!)$$

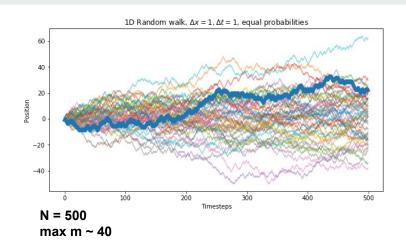


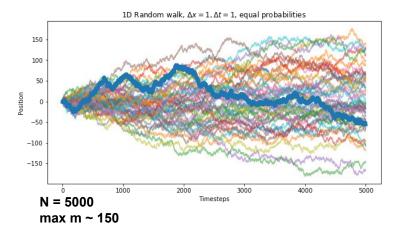
Simulating a random walk



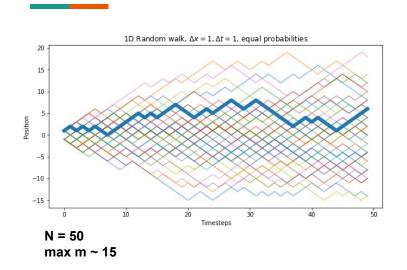
Plotted with p = 0.5

Note that the fluctuations have the same "structure" at any N - relative size of the variations in the path



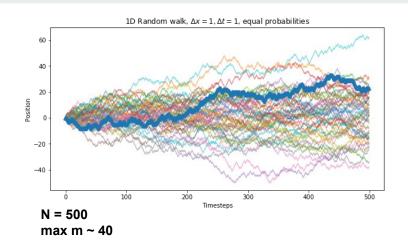


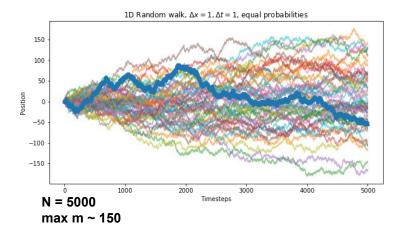
Simulating a random walk



Note that the spread of the random walks increases as roughly the square root of N

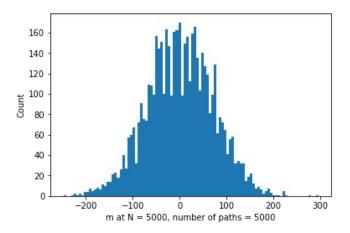
This is a consequence of the Central Limit Theorem $\rightarrow \sigma$ proportional to \sqrt{N}





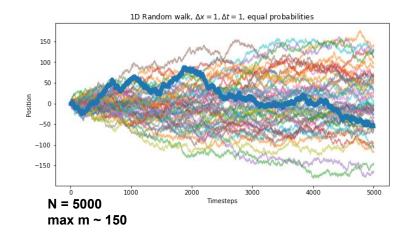
Simulating a random walk

Gaussian distribution of end points (5000 trials, N = 5000 timesteps)



Note that the spread of the random walks increases as roughly the square root of N

This is a consequence of the Central Limit Theorem $\rightarrow \sigma$ proportional to \sqrt{N}



Taking the continuous limit for a 1-D walk

- Use a non-drifting walk, i.e. $p = \frac{1}{2}$
- Probability of arriving at some (m, N) is $\frac{1}{2}$ the probability of arriving at $(m-1,N-1) + \frac{1}{2}$ the probability of arriving at (m+1,N-1), at the previous timestep
 - o i.e. $w(m,N) = \frac{1}{2} w(m-1,N-1) + \frac{1}{2} w(m+1,N-1)$
- Switching from m,N to x,t: $m = x/\Delta x$ and $N = t/\Delta t$. Let u = continuous probability function.
 - $\circ \qquad 2u(x,t) = u(x-\Delta x, t-\Delta t) + u(x+\Delta x, t-\Delta t)$

Taking the continuous limit for a 1-D walk

• Take the Taylor expansion around x and t with the variables Δx and Δt .

$$2u = u - \Delta t^* u_t - \Delta x^* u_x + \frac{1}{2} (\Delta x^2 u_{xx} + 2\Delta x \Delta t u_{xt} + \Delta t^2 u_{tt}) + ...$$

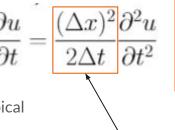
$$+ u - \Delta t^* u_t + \Delta x^* u_x + \frac{1}{2} (\Delta x^2 u_{xx} - 2\Delta x \Delta t u_{xt} + \Delta t^2 u_{tt}) + ...$$

$$\circ \qquad 0 = 2 \, \Delta t^* u_t^1 + \Delta x^2 \, u_{xx}^2 + \Delta t^2 \, u_{tt}^2 + \dots$$

• At small Δt and Δx , we get the following:

○
$$u_{t} = (\Delta x^{2}/2\Delta t) u_{yy}$$
, also stated \rightarrow

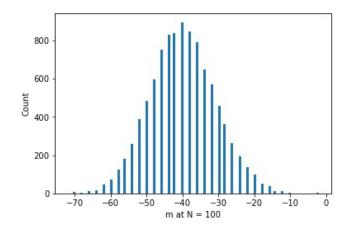
- Constant is the diffusion coefficient D
- → Higher D → faster diffusion due to larger typical spacial step in a random walk
- Only linearly varying distributions are stable
- In 1-D, solutions look like linear connections between any boundary conditions

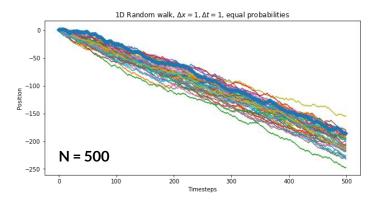


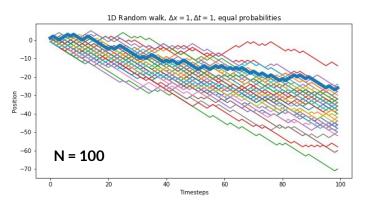
The diffusion equation describes changes in distributions where motion is governed by random walks

Other walks - drifting 1-D walk

- Unequal p and q
- Final positions still obey Gaussian distribution:



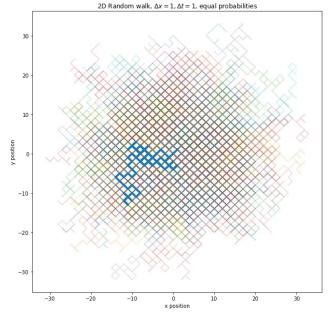


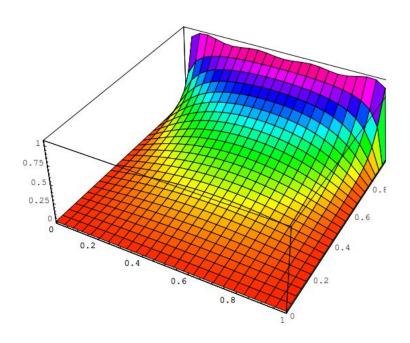


p = 0.3, and q = 0.7

- Different probabilities for (up, left), (up, right), (down, left), (down, right)
- For equal probabilities, is essentially two random walks in different directions superimposed
- Can make it drift diagonally, or create some correlation between different choices for up/down and right/left
- Creates the diffusion equation as follows:
- Rate of change depends on curvature - configurations that solve Laplace's equation are stable

$$rac{\partial \phi({f r},t)}{\partial t} = D
abla^2 \phi({f r},t),$$





Solutions for various boundary conditions can be solved with Laplace's equation

